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Key Points:

- Akatsuki/LIR revealed the global structures of thermal tides across the equator in the upper cloud layer of Venus for the first time
- Using the emission angle dependence of LIR's sensing altitude, upward propagation of the semidiurnal tide was confirmed
- Wave types consisting of the thermal tides were identified

Supporting Information:

- Supporting Information S1

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Global Structure of Thermal Tides in the Upper Cloud Layer of Venus Revealed by LIR on Board Akatsuki

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Abstract Longwave Infrared Camera (LIR) on board Akatsuki first revealed the global structure of the thermal tides in the upper cloud layer of Venus. The data were acquired over three Venusian years, and the analysis was done over the areas from the equator to the midlatitudes in both hemispheres and over the whole local time. Thermal tides at two vertical levels were analyzed by comparing data at two different emission angles. Dynamical wave modes consisting of tides were identified; the diurnal tide consisted mainly of Rossby-wave and gravity-wave modes, while the semidiurnal tide predominantly consisted of a gravity-wave mode. The revealed vertical structures were roughly consistent with the above wave modes, but some discrepancy remained if the waves were supposed to be monochromatic. In turn, the heating profile that excites the tidal waves can be constrained to match this discrepancy, which would greatly advance the understanding of the Venusian atmosphere.

Plain Language Summary On Venus, the atmosphere circulates 60 times faster than the solid body of Venus; this phenomenon is called “superrotation,” and it is one of the mysteries of the Venusian atmosphere. To maintain the fast circulation, thermal tides, which are global-scale atmospheric waves excited by solar heating, have been considered a very important candidate because they have the ability of accelerating the atmosphere through propagating. A midinfrared camera onboard the Japanese Venus orbiter, Akatsuki, can capture temperature perturbations due to the thermal tides in the upper cloud level (60- to 70-km altitude), and it revealed their global and vertical structures with a long-term observation (more than three Venusian years) for the first time. Interestingly, we found that the location of the maximum temperature at the cloud top level was different from noon where solar energy input is at a maximum. In addition, the location was shifted toward the morning side as the sensing altitude increased. This finding is an evidence of the vertical traveling of the thermal tides, indicating the wave's atmospheric acceleration.

1. Introduction

Thermal tides excited in the cloud layer has been considered one of the main acceleration sources that maintain the atmospheric superrotation of Venus (e.g., Fels & Lindzen, 1974; Hou et al., 1990; Plumb, 1975; Newman & Leovy, 1992; Takagi & Matsuda, 2007), in which the zonal wind speed of the atmosphere at the cloud top (~70 km) is more than 60 times faster than the rotation speed of the solid body of Venus. The structures of the thermal tides have been confirmed in a temperature field from space- and ground-based observations (e.g., Ainsworth & Herman, 1978; Apt et al., 1980; Migliorini et al., 2012; Taylor et al., 1980; Zasova et al., 2007) and in zonal and meridional wind fields at the cloud top level by tracking cloud motions (e.g., Horinouchi et al., 2018; Kouyama et al., 2012; Limaye & Suomi, 1981; Moissl et al., 2009; Rossow et al., 1990; Sánchez-Lavega et al., 2008). However, despite the importance of thermal tides in the Venusian atmosphere, the global feature of thermal tides across all local times and latitudes has not been obtained due to the limited data coverage of previous observations (e.g., a ground-based observation did not cover the subsolar region, Apt et al., 1980; cloud tracking was limited to only dayside, and only one hemisphere was observable by spacecraft with a polar orbit, Migliorini et al., 2012; Taylor et al., 1980). Ground-

based observations by Ainsworth and Herman (1978) covered the northern and southern hemispheres simultaneously, but the observational dates were quite limited, resulting in large data gaps in the high latitudes. Since the tidal amplitude at middle to high latitudes from their results is much greater than those obtained by space-borne observations, verification is needed.

In this study, latitudinal profiles of diurnal, semidiurnal, and higher-frequency components of thermal tides in the temperature field were investigated based on long-term observation data from Longwave Infrared Camera (LIR) on board Akatsuki (Taguchi et al., 2007). LIR detects thermal emission from the Venusian upper cloud layer (60–70 km) and has been observing Venus continuously since December 2015. Thanks to the Akatsuki's equatorial orbit, the whole globe of Venus can be observed. These sufficiently long-term and global observations allow us to study the global structure of thermal tides over the whole local time and for a wide latitudinal range in both hemispheres reliably. Because latitudinal phase variations in the thermal tides indicate horizontal momentum transportation, the global profile of the thermal tides is crucial for understanding its impact in the Venusian atmosphere. In addition to the horizontal structure, this study accesses the vertical structure of the thermal tides through a comparison of the tidal structures at different altitudes, using the emission angle dependence of sensing altitudes. To the best of our knowledge, this paper is the first to discuss the wave type of the thermal tides based on the long-term observational results.

2. Observations and Data

LIR covers wavelengths of 8–12 μm and captures thermal emissions from the cloud top level of Venus (Taguchi et al., 2007). An advantage of LIR is that it can observe both the dayside and nightside of Venus. LIR continuously observes Venus at 1- or 2-hr intervals, resulting in 8–15 observations every Earth day except during the superior conjunction, where the spacecraft cannot communicate with the tracking station on Earth for ~ 1 month. In this study, we used more than 22,000 LIR images obtained from October 2016 to January 2019 (more than three Venusian years). The spatial resolution of the images of Venus changes according to the distance between Venus and Akatsuki. Near the periapsis, the resolution is better than 10 km/pixel, whereas it reaches more than 300 km/pixel at the apoapsis altitude of $\sim 360,000$ km where the radius of the Venus disk is 20 pixels. The LIR images used in this study were obtained from the AKATSUKI Science Data Archive (Murakami et al., 2017).

The noise equivalent temperature difference is 0.3 K at the target temperature of 230 K, which corresponds to the temperature resolution of LIR. The absolute temperature uncertainty is 3 K (Fukuhara et al., 2011). The primary cause of the absolute temperature error is variations of instrument's thermal condition affected by the attitude of the spacecraft. The variation due to the thermal condition was almost equally distributed across all local solar time regions, and thus, the effect of the absolute error on the brightness temperature map after averaging (shown later) is expected to be much smaller than the original absolute error of 3 K. The relative uncertainty from the noise equivalent temperature difference is also reduced by the averaging. In total, the temperature uncertainty for tidal components is estimated to be 0.1 K in the present analysis. The LIR images have a background bias that strongly depends on the baffle temperature of LIR. This bias was determined from deep space observations and successfully removed in all LIR image (Fukuhara et al., 2017). An additional temperature calibration was performed to reduce the unexpected gradual temperature increase seen in deep space images (supporting information Text S1).

The contribution function representing the sensing altitude of LIR is centered at 65-km altitude with a full width at half maximum of ~ 10 km for a nadir viewing geometry (Figure 1). The contribution function was calculated by a line-by-line radiative transfer calculation following Sato et al. (2014) with the spectral response of LIR (Taguchi et al., 2007). The calculation adopted the equatorial ($<30^\circ$) temperature profile from the Venus International Reference Atmosphere (Seiff et al., 1985), the gaseous absorption profiles from Marcq et al. (2005), and the cloud particle number density from Haus et al. (2014), which assumes four different-sized particle modes (mode 1, mode 2, mode 2', and mode 3), composed of 75% H_2SO_4 and 25% H_2O . Then, the line-by-line results were convolved with the filter transmittance of LIR (Taguchi et al., 2007)

LIR observes higher altitudes for larger emission angles (e.g., Taguchi et al., 2012; supporting information Text S2). According to the vertical temperature profile of the Venusian atmosphere (negative lapse rate), the observed brightness temperature becomes cooler at higher emission angles (i.e., limb darkening) in the low and middle latitudes. In other words, the vertical structures of the thermal tides can be

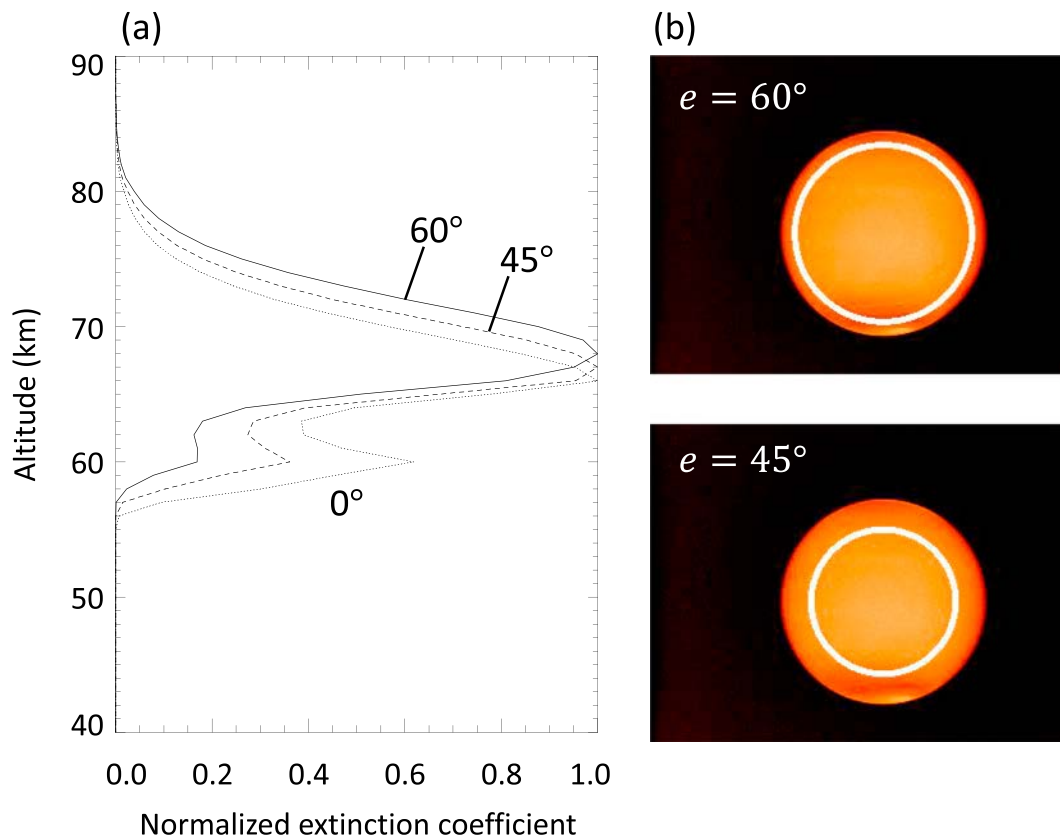


Figure 1. (a) Contribution functions of the three emission angles, 60° (solid), 45° (dashed), and nadir (dotted), based on the Longwave Infrared Camera's (LIR's) spectral response. (b) Examples of data coverage in an LIR image for emission angles of 60° and 45°, whose positions are indicated by white circles.

investigated by comparing brightness temperature maps from different emission angles, although the wide width of the LIR's contribution function does not allow precise altitude assignments. To investigate the vertical structures of the thermal tides, we selected two specified emission angles: 60° ($\pm 3.5^\circ$) to obtain a wide latitudinal coverage and 45° ($\pm 2.5^\circ$) for comparison (Figure 1). The effective altitudes of the contribution function are 68.9 and 67.6 km for the emission angles of 60° and 45°, respectively, which are evaluated from

$$z_e = \frac{\int z f_e(z) dz}{\int f_e(z) dz}, \quad (1)$$

where z denotes the altitude and f_e represents the LIR's contribution function for each emission angle, e . This indicates a 1.3-km difference in the sensing altitudes between the two emission angles. The uncertainty of the altitude difference can be less than 0.3 km, which was confirmed by the results comparing with a temperature profile obtained from Akatsuki radio occultation observations (Imamura et al., 2017; supporting information Text S2). To obtain a continuous data set, since the spacecraft was not always above the equator and the latitudinal coverage changes from observation to observation, the latitudinal range of $\pm 55^\circ$ was analyzed for the emission angle of 60°, while $\pm 40^\circ$ for 45°.

3. Global Structure of Thermal Tides

To focus on the solar-fixed component in the temperature field, we first averaged the observed brightness temperature as a function of the local time and latitude across the whole observation period. Then, the temperature perturbation, which should mainly consist of thermal tides, was extracted by subtracting the zonal mean brightness temperature at each of the 5° latitudinal bins. Figure 2a shows the global structure of the

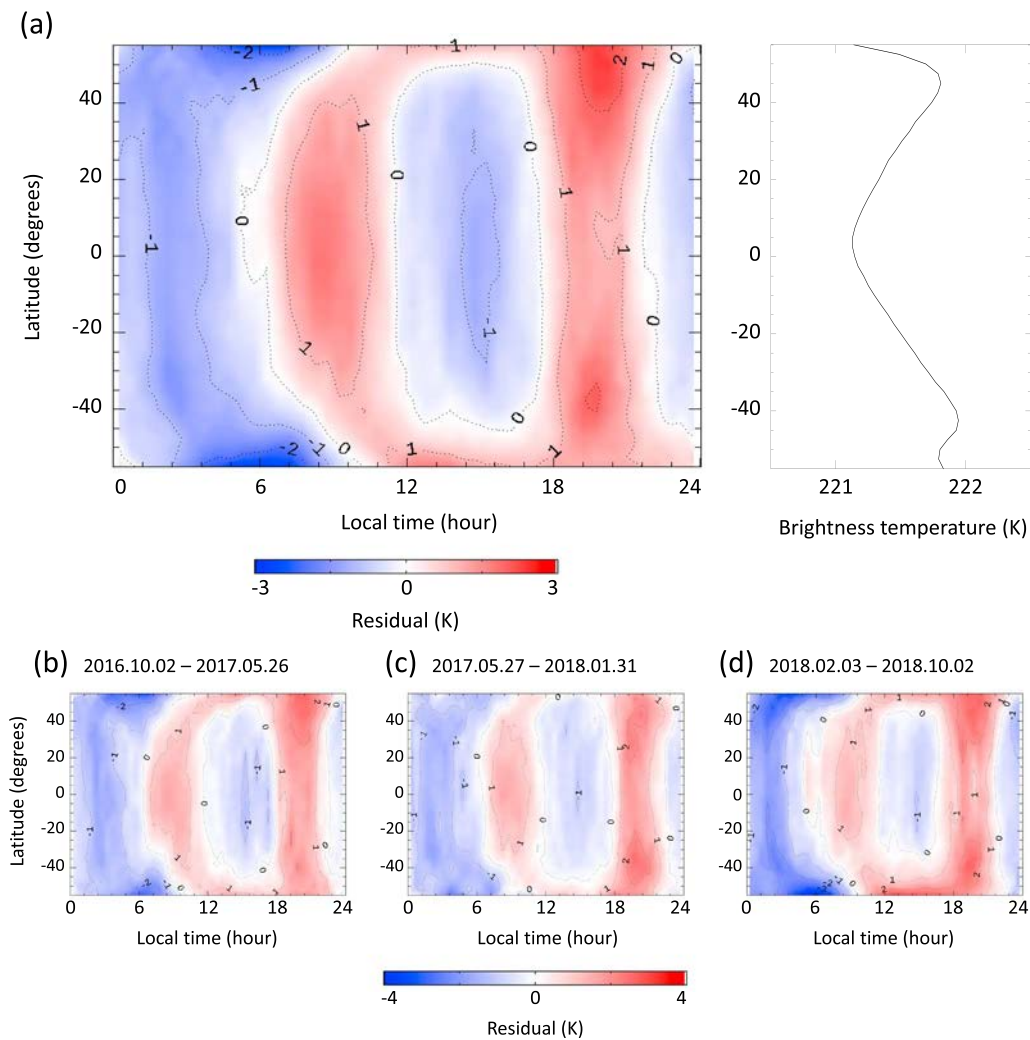


Figure 2. (a) Thermal tide structure obtained with a 60° emission angle condition by using the whole data period from October 2016 to January 2019. The brightness temperature in the Longwave Infrared Camera (LIR) images was averaged for every 0.5 hr local time bin and 5° latitudinal bin, and then the zonal mean temperature was subtracted (right panel of [a]) at each latitude to extract the temperature perturbation. The typical standard error for a grid was 0.08 K. (b–d) Same as (a) but for each Venusian year over three Venusian years.

thermal tides without any local time gaps, which was obtained for the first time in the history of Venus observations. Equatorially symmetric structures clearly appeared at all local times. Notably, a local temperature maximum was observed at approximately 9 hr in the equatorial region, not at the local noon where solar incident energy is at a maximum. Local peaks at approximately 19 hr were observed at the midlatitudes of both hemispheres. It is worth mentioning that the tidal structure was almost the same for every Venusian year, especially at latitudes lower than 40° (Figures 2b–2d). The structures shown in Figure 2a can be considered typical structures of the thermal tides in the Venusian atmosphere during, at least, the Akatsuki observation period.

It should be noted that sensed altitude perturbation can also affect the observed temperature in addition to the atmospheric temperature perturbation, and they cannot be distinguished by the LIR observation alone. In this study we assumed that the temperature fields in Figure 2 were from only atmospheric temperature perturbation, since an effect from the altitude perturbation due to a tidal component can be reduced by the integration effect of the LIR contribution function (the magnitude can be 0.1 K; supporting information Text S3). This assumption should be validated by a combination of observations for different altitudes, such as radio occultation and midinfrared spectral observations in future works.

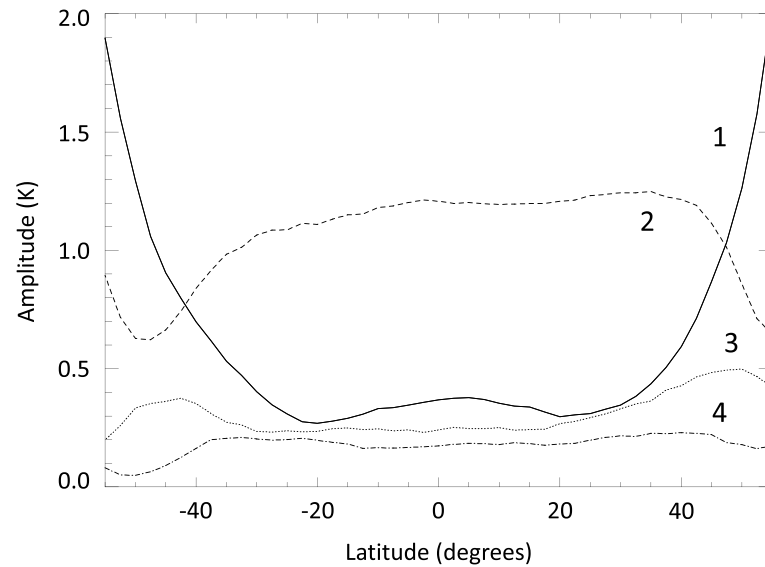


Figure 3. Latitudinal amplitude profiles of diurnal (solid line), semidiurnal (dashed), terdiurnal (dotted), and quarter-diurnal (dash-dotted) tides in the 60° emission-angle data set.

To clarify which tidal component contributed to the average structure, latitudinal profiles of amplitudes for wavenumber 1–4 components, namely, the diurnal, semidiurnal, terdiurnal, and quarter-diurnal tides, were investigated by fitting sinusoidal functions of wavenumbers 1–4 as

$$f = A_1 \sin(x + \varphi_1) + A_2 \sin(2x + \varphi_2) + A_3 \sin(3x + \varphi_3) + A_4 \sin(4x + \varphi_4) \quad (2)$$

with a nonlinear least squares method (Figure 3), where x denotes local time and A_1 – A_4 and φ_1 – φ_4 are fitted parameters. The error in each amplitude was less than 0.05 K, evaluated from the fitting. In the low latitudes, the semidiurnal tide was the most significant component with an amplitude greater than 1 K, and its amplitude remained to ~40° latitude. On the other hand, the diurnal tide became significant at middle to higher latitudes where the amplitude of the semidiurnal tide decreased. These tendencies were the same as those at the cloud top level derived from ground-based observations in 1977–1979 (Apt et al., 1980) and the Pioneer Venus Orbiter Infrared Radiometer observation in the northern hemisphere (Taylor et al., 1980), indicating long-term steadiness of the thermal tide structure. Note that the amplitude observed by LIR were almost a half of those in the Orbiter Infrared Radiometer observations. In addition to possible temporal variations in the amplitudes, the difference could be from LIR's long-term observations, which may be more suitable to catch mean structures than the previous observation (10 weeks); however, the LIR's wider contribution function may also affect the amplitude, which integrates the vertical structure of a vertically tilting tidal component and thus smoothens the observed tidal structure. From a simple test for evaluating the integration effect, the actual amplitude of the semidiurnal tide could be twice as large as the observed amplitude (supporting information Text S3).

The terdiurnal tide showed small peaks in the midlatitudes of both hemispheres, which is consistent with a global circulation model (GCM; Takagi et al., 2018). Whole components showed almost hemispherical symmetry, but the amplitude of the semidiurnal tide in the northern hemisphere was slightly stronger than that in the southern hemisphere. Recently, an asymmetric structure was reported in zonal winds at the cloud top (Horinouchi et al., 2018). There could be a mechanism that induces the asymmetric condition in the Venusian atmosphere whose impact is worth investigating with a numerical approach.

The analysis of the tidal components also provides latitudinal profiles of phases from which a horizontal structure of each tidal component can be reproduced (Figure 4). The diurnal tide showed a clear latitudinal phase tilt toward the evening (westward) direction from the equator to the midlatitudes. The semidiurnal tide had a flat phase profile from the low to middle latitudes, which was in good agreement with the

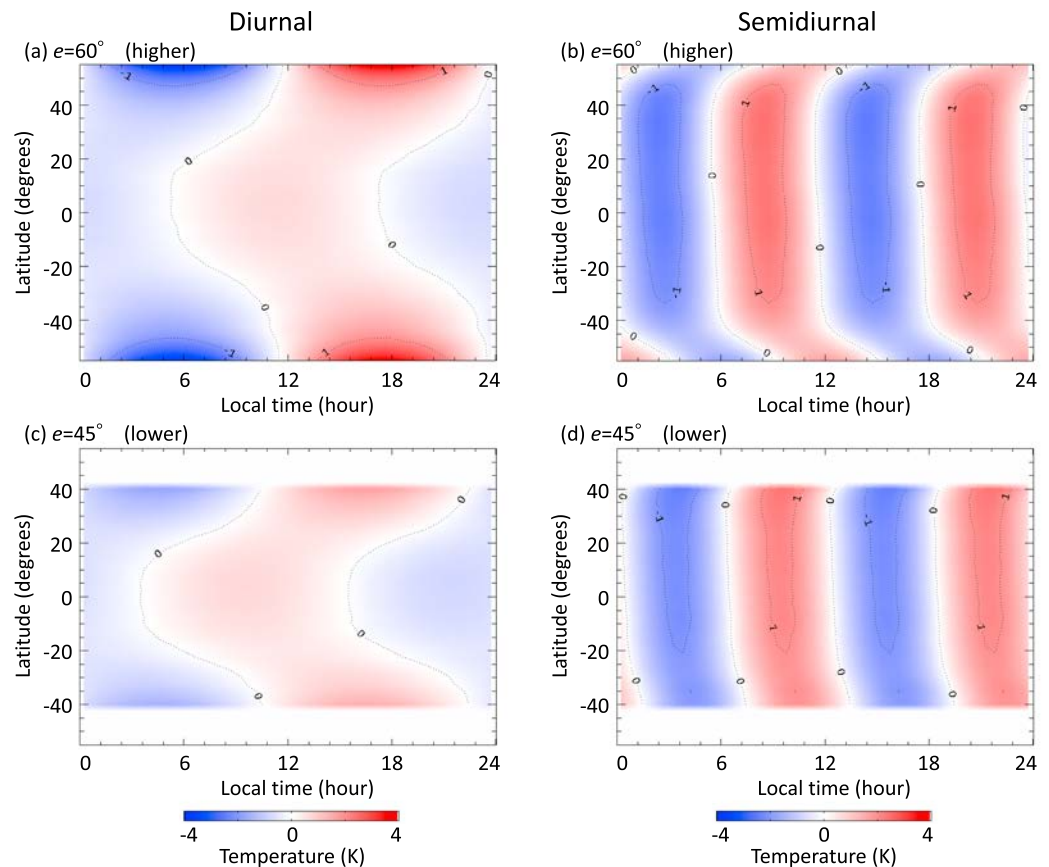


Figure 4. Horizontal structures of (a) diurnal and (b) semidiurnal components obtained with the 60° emission angle (higher altitudes) and (c, d) those with the 45° emission angle (lower altitudes). Note that the direction of the superrotation is from morning to evening.

latitudinal range of the significant amplitude. On the other hand, the semidiurnal tide showed a clear phase tilt toward the evening direction around the midlatitudes.

The local maximum of the semidiurnal tide clearly shifted from noon, whereas the diurnal tide had a local maximum approximately at noon in the equatorial region. Since solar heating is at a maximum at noon and it is uniformly zero on the night side, the components of both wavenumbers 1 and 2 of the heating should have their local maxima at noon at their excitation altitude. Therefore, the phase shift observed in the semidiurnal tide indicates vertical propagation of the wave from its excitation altitude, although the excitation altitude has not yet been determined.

4. Discussion

4.1. Wave Modes of Thermal Tides and Their Implications

The classical tidal theory based on Laplace's tidal equation (Longuet-Higgins, 1968) is useful to interpret the observed tidal waves in this study. The superrotating zonal winds act as the background rotation for the thermal tides. If we assume solid-body background rotation, the nondimensional parameters σ and γ introduced by Longuet-Higgins (1968) have the same values for each of the diurnal or the semidiurnal tide, irrespective of the planet (whether on Venus or on Earth); see the supporting information (Text S4) for the definition of the parameters and the evaluation of the dimensional parameters. The numerical study by Takagi and Matsuda (2005) suggests that Venusian tidal waves on realistic background winds are somewhat similar to those on the uniform rotation, so the actual tidal waves can be interpreted as modifications of the tidal modes in the theory of Longuet-Higgins (1968).

From the tidal theory, the diurnal tidal features in Figures 4a and 4c are interpreted as the superposition of (mainly) two modes: the gravest symmetric Rossby-wave mode with a negative equivalent depth and the gravest symmetric gravity-wave mode, as is the case for the atmospheric thermal tides on Earth (e.g., Chapman & Lindzen, 1970). The diurnal Rossby-wave mode has a large amplitude at the middle to high latitudes. The mode is vertically evanescent, so the temperature disturbance should be nearly in phase with the geopotential height disturbance. In fact, the observed cloud top meridional winds of thermal tides (Horinouchi et al., 2018) are consistent with this temperature-geopotential relation. The vertical evanescence is also indicated from the tidal structure in Figures 4a and 4c (see section 4.2). The diurnal gravity-wave mode has an equivalent depth of ~ 40 m, corresponding to a vertical wavelength of 6 km (supporting information Text S4). Such a short vertical wavelength was confirmed in a linearized calculation (Pechmann & Ingersoll, 1984). The corresponding equatorial radius of deformation is $\sim 2.5 \times 10^3$ km, so the mode has nodes of temperature and geopotential height at $25\text{--}30^\circ\text{N/S}$. Then, the amplitude at higher latitudes should be quite small that it was inconsistent from the observation, which is another evidence that the diurnal tide cannot be interpreted solely as a gravity wave.

The observed structure of the diurnal tide is possibly superposition of the two wave modes, and such superposition may cause a complex phase relationship between temperature and geopotential height. Therefore, it is difficult to assess momentum transport due to the diurnal tide from the observed latitudinal phase variation of the diurnal component even qualitatively. Moreover, the short vertical wavelength of the gravity-wave mode introduces an additional complexity for evaluating momentum transportation due to the diurnal tide discussed in section 4.2. However, one can expect that the diurnal tide may carry the westward (along with the superrotation) momentum toward low latitudes since its thermal forcing mainly resides at low latitudes.

The solution for the semidiurnal tide in the tidal theory consists only of gravity waves. The semidiurnal temperature disturbances in Figures 4b and 4d indicate the gravest symmetric gravity-wave mode. This mode is expected to have a vertical wavelength of 22 km, and the corresponding radius of equatorial deformation is over 40° (supporting information Text S4), which is consistent with the observed flat structure of the semidiurnal tide. It is rather difficult to interpret the phase tilt at midlatitudes in Figure 4b, because it may reflect a gradual decrease in the cloud top level from middle to high latitudes (Ignatiev et al., 2009) and a midlatitude jet structure.

4.2. Vertical Propagation of the Tides

The emission angle dependence of the effective altitude enables us to investigate the vertical structure of the tidal waves. In the diurnal tide (Figures 4a and 4c), there is no apparent phase difference at the midlatitudes between the emission angles of 45° and 60° . This feature is consistent with the vertical evanescence of the Rossby-wave mode. At low latitudes, there exists a phase difference consistent with the westward tilt (toward the evening side) with altitude. This result indicates downward energy (group) propagation. However, since the depth of the solar heating should be much broader than half of the theoretical vertical wavelength, 6 km, a strong superposition is expected to arise from the vertical heating structure. Therefore, it is difficult to conclude whether the wave is truly propagating downward from the cloud top altitude. Nevertheless, if we assume an altitude difference of 1.3 km between the two sensing altitudes (higher for the emission angle of 60° and lower for the emission angle of 45°) and if we tentatively assume a vertically monochromatic wave near the equator, the phase difference indicates that the vertical wavelength is ~ 10 km. Note that the vertical superposition owing to the heating distribution is the likely cause of the vertical evanescence of the simulated diurnal tide around the cloud top level even near the equator according to the GCM (Takagi et al., 2018).

On the other hand, the phase difference of the semidiurnal tide had the eastward tilt (toward the morning side) with altitude, which is consistent with the direction of upward energy propagation. Quantitatively, the phase difference was 0.86 hr, which indicates that the vertical wavelength of the semidiurnal tide was 18 ± 4 km if we assume a 1.3 ± 0.3 -km altitude difference between the two sensed altitudes, and this wavelength was almost the same as the estimation from radio occultation observations (Ando et al., 2018). In addition, the measured vertical wavelength is also consistent with the theoretical expectation of 22 km. Therefore, the actual group propagation is likely upward. Further study is needed to elucidate the full emission angle dependency of the thermal tide phases.

Vertical propagation of the semidiurnal tide at the cloud level was reproduced in several GCMs (Lebonnois et al., 2016; Takagi et al., 2018; Yamamoto et al., 2019). However, the phases of the modeled semidiurnal tides at the cloud level were different in different models, and some model results were different from the LIR result. For example, Takagi et al. (2018) showed that the semidiurnal tide has a local temperature maximum of approximately 15 hr at the cloud level in their model, whereas our result suggests the local temperature maxima around 9 and 21 hr. Considering the vertical propagation of the semidiurnal tide, the difference may provide a constraint for the vertical profile of the solar heating in a model, because the wave phase at an altitude should be sensitive to the excitation altitude of the wave that is affected from the solar heating profile.

5. Conclusions

In this study, we presented a global structure of thermal tides in the upper cloud layer of Venus using LIR's long-term data, acquired from Akatsuki's equatorial orbit. This global structure was revealed for the first time in the history of Venus ground- and space-based observations. The observed structure of the thermal tides showed a good consistency with previous observations, indicating the steadiness of the tides in the Venusian atmosphere. There were fewer variations in the thermal tide structure during the three Venusian years analyzed in this study.

Tidal components were investigated with an analysis for periodicity. The extracted structure of the diurnal tide, which had a small amplitude at low latitudes and a larger amplitude at higher latitudes, indicated a superposition of the gravest symmetric Rossby-wave and gravity-wave modes. On the other hand, the semidiurnal tide component was significant in the lower to middle latitudes, and its flat phase structure was consistent with the gravest symmetric gravity-wave mode. The phase difference in the semidiurnal structures between the emission angles of 45° and 60° was consistent with upward energy propagation and a vertical wavelength of ~18 km, which is close to the theoretically expected value. A similar emission angle comparison for the diurnal tide suggested vertical evanescence in the midlatitudes, which is also consistent with the tidal theory.

Because recent observations of Venus suggest a global albedo variation for Venus (Lee et al., 2015, 2019) that should affect the excitation of thermal tides, longer-term LIR data will be useful in monitoring possible temporal variation in the thermal tides. Since LIR observations can allow to retrieve the vertical structure of the atmosphere by utilizing the emission angle dependence of sensing altitude, a combination of such LIR observations and radio occultation observations (cf. Ando et al., 2018) will clarify the excitation altitude of the tides and their vertical propagations. By combining the LIR results with numerical modeling, it is also possible to deduce the vertical structure of heating that excites thermal tides, which would greatly advance our understanding of the dynamics of the Venusian atmosphere. In addition, the thermal tide structure derived in this study may provide a more realistic three-dimensional atmospheric condition that will help to discuss a local temperature profile for small-scale convection activity (cf. Lefevre et al., 2018).

Acknowledgments

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